

Low palaeopressure of the martian atmosphere estimated from the size distribution of ancient craters

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The decay of the martian atmosphere—which is dominated by carbon dioxide—is a component of the long-term environmental change on Mars¹ from a climate that once allowed rivers to flow^{2–6} to the cold and dry conditions of today. The minimum size of craters serves as a proxy for palaeopressure of planetary atmospheres, because thinner atmospheres permit smaller objects to reach the surface at high velocities and form craters^{7–9}. The Aeolis Dorsa region near Gale crater on Mars contains a high density of preserved ancient craters interbedded with river deposits¹¹ and thus can provide constraints on atmospheric density at the time of fluvial activity. Here we use high-resolution images and digital terrain models¹⁰ from the Mars Reconnaissance Orbiter to identify ancient craters in deposits in Aeolis Dorsa that date to about 3.6 Gyr ago and compare their size distribution with models of atmospheric filtering of impactors^{12,13}. We obtain an upper limit of 0.9 ± 0.1 bar for the martian atmospheric palaeopressure, rising to 1.9 ± 0.2 bar if rimmed circular mesas—interpreted to be erosionally-resistant fills or floors of impact craters—are excluded. We assume target properties appropriate for desert alluvium¹⁴: if sediment had rock-mass strength similar to bedrock at the time of impact, the paleopressure upper limit increases by a factor of up to two. If Mars did not have a stable multibar atmosphere at the time that the rivers were flowing—as suggested by our results—then a warm and wet CO₂/H₂O greenhouse² is ruled out, and long-term average temperatures were most likely below freezing.

Planetary atmospheres brake, ablate and fragment small asteroids and comets, filtering out small high-velocity surface impacts and causing fireballs, airbursts, meteors and meteorites. The smallest impact craters near sea level on Earth have diameter $D \sim 20$ m. ‘Zap pits’ as small as $30 \mu\text{m}$ are known from the airless Moon, but other worlds show the effects of progressively thicker atmospheres: the modern martian atmosphere can remove $>90\%$ of the kinetic energy of >240 kg impactors⁷; Titan’s paucity of small craters is consistent with atmospheric filtering of craters smaller than 6–8 km (ref. 8); and on Venus, craters $D < 20$ km are substantially depleted by atmospheric effects⁹.

Changes in the concentration of atmospheric volatiles are believed to be the single most important control on the climate evolution and habitability of Mars, which in turn is a benchmark for habitable-zone calculations for exoplanets¹⁵. Contrary to early work², it is doubtful that increasing CO₂ pressure (\approx total atmospheric pressure, P) is enough to raise the mean annual surface

temperature (\bar{T}) of early Mars to the freezing point, even when water vapour and cloud feedbacks are considered⁵. However, increased CO₂ aids transient surface liquid water production by impacts, volcanism, or infrequent orbital conditions^{3,4,6}. Existing data require an early epoch of massive atmospheric loss to space, suggest that the present-day rate of escape to space is small, and offer evidence for only limited carbonate formation¹⁶. These data have not led to convergence among atmosphere evolution models, which must balance poorly understood fluxes from volcanic degassing, escape to space, weathering and photolysis¹⁷. More direct measurements¹⁸ are required to determine the history of Mars’s atmosphere. Wind erosion exposes ancient cratered volumes on Mars and the size of exhumed craters has been previously suggested as a proxy of ancient Mars P (for example, ref. 19).

Here we obtain a new upper limit on the atmospheric pressure of early Mars from the size–frequency distribution of small ancient craters (Fig. 1) interspersed with river deposits in Aeolis, validated using High Resolution Imaging Science Experiment (HiRISE) digital terrain models (DTMs) and anaglyphs, in combination with simulations of the effect of P on the crater flux. The craters are interbedded with river deposits up to $\sim 10^3$ km long, with inferred peak river discharge $10\text{--}1,000 \text{ m}^3 \text{ s}^{-1}$ (ref. 10). Therefore, the atmospheric state they record corresponds to an interval of time when Mars was substantially wetter than the present, probably >3.6 Gyr ago (Ga) (Supplementary Methods).

Aeolis Dorsa (part of the Medusae Fossae Formation) is a promising location to hunt for ancient craters: the stratigraphy contains numerous channel-fill deposits of large rivers and, when these overlie a crater, that crater must be as ancient as the rivers²⁰. Certain beds in Aeolis Dorsa (Supplementary Methods) preserve a high density of ancient craters, perhaps due to slow deposition or a diagenetic history unusually favourable for crater preservation. We constructed stereo DTMs/orthophotos for two image pairs covering these beds, DTM1 and DTM2 (Methods). Following a checklist (Supplementary Table), craters were classified as definite ancient craters (visibly embedded within stratigraphy: for example, overlain by river deposit; $n = 56$, median diameter $D_{50} = 107$ m, 10th-percentile diameter $D_{10} = 50$ m), rimmed circular mesas (RCMs; $n = 71$, $D_{50} = 48$ m, $D_{10} = 21$ m), or candidate ancient crater ($n = 192$, D_{50} also 48 m, D_{10} also 21 m; candidates are not considered further, but their inclusion would strengthen our conclusions). We measured D by fitting circles to preserved edges/rims. RCMs appear as disks in raw HiRISE images. We interpret them as the erosionally resistant fills/floors of impact craters that were topographically inverted during the deflation of the target unit. They are unlikely to be

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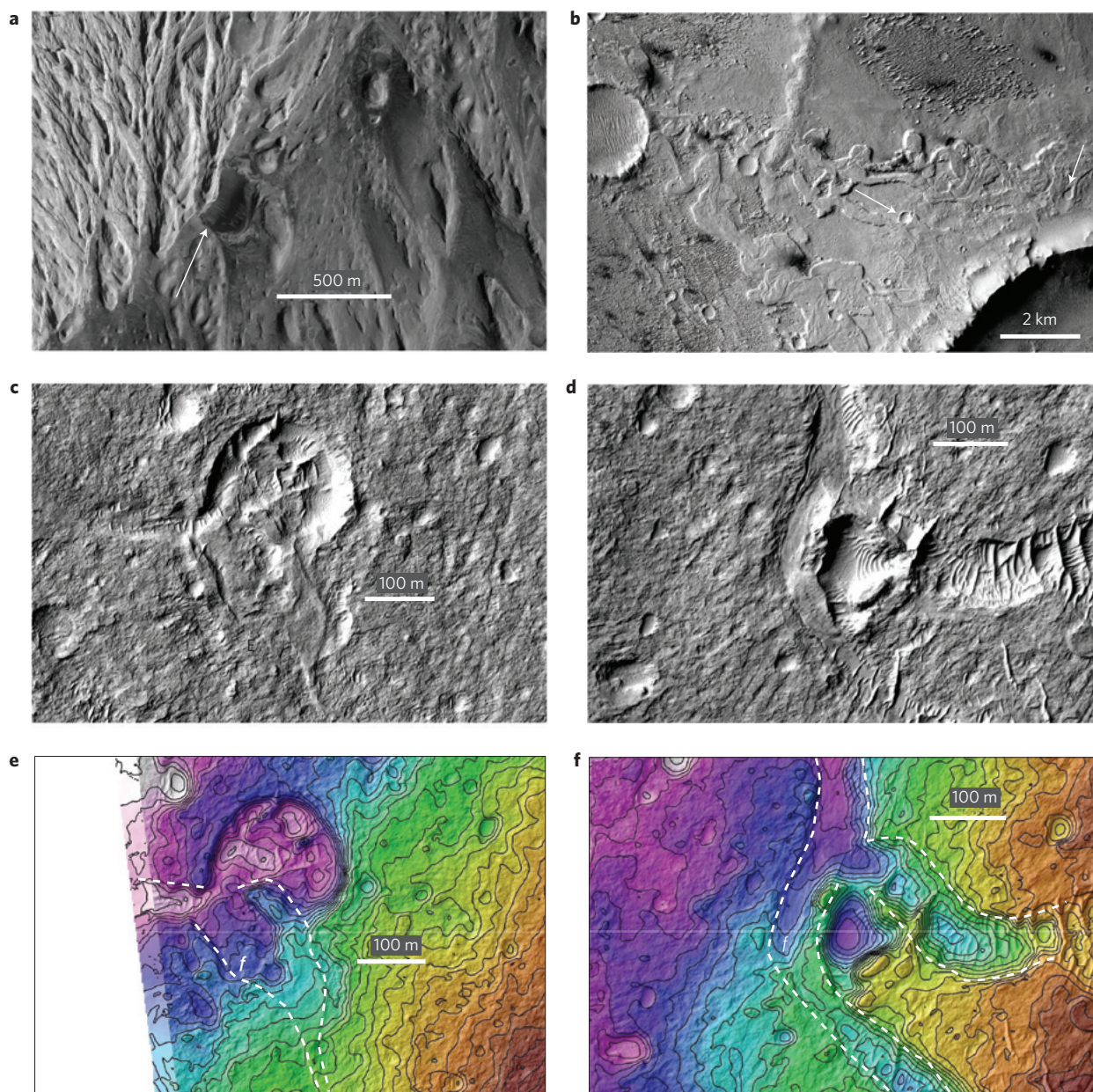


Figure 1 | Gallery of ancient martian craters (from ref. 20). **a**, Crater being exhumed from beneath an unconformity within Gale crater's mound (Aeolis Mons/Mount Sharp), the destination of the Curiosity rover. ESP_019988_1750. **b**, Craters with intact rims being exhumed from beneath meander belts, Aeolis Dorsa, G03_019249_1744_XI_05S205W; ref. 10. **c**, Crater partly draped by fluvial channel materials (f in **e**), Aeolis Dorsa. Diameter, 238 m. ESP_019104_1740. **d**, Crater partly draped by fluvial channel materials (f in **f**), Aeolis Dorsa, ESP_019104_1740. Diameter, 141 m. **e**, Crater from **c**, but with 1 m elevation contours from DTM2 (see text). **f**, Crater from **d** with 1 m contours from DTM2.

outliers of a young mantle because they are not found away from the fluvial unit. We plot them separately, but consider them to be probable ancient craters. We used unambiguously ancient craters as a guide to the preservation state of the smaller craters. These ancient craters are unlikely to be maars; maars are not randomly distributed in space or time/stratigraphy. We also reject the possibility that they are palaeokarst sinkholes; sinkholes lack rims, are concentrated at particular stratigraphic levels, and are overdispersed.

We generated synthetic crater populations for varying P (ref. 12). The approach is conceptually similar to that of previous studies¹³ and benefits from measurements of the present martian cratering flux (Methods). Modelled smallest-crater diameter increases linearly with pressure (~ 20 m at 1 bar) as expected from equating impactor and atmospheric-column masses²¹. This is broadly

consistent with low-elevation impacts on Earth (the column mass of Earth's sea-level atmosphere is equivalent to ~ 0.4 bar on Mars). We apply a geometric correction for exhumation from a cratered volume (Supplementary Methods) assuming that initial crater shape is isometric over the diameter range. After bayesian fitting, we correct our P estimate for elevation (our DTMs are below datum; Mars's average P was 20% lower than local P).

We compared the model with the combined data set (DTM1 + DTM2). Combined best fits are $P = 1.9 \pm 0.2$ bar, falling to $P = 0.9 \pm 0.1$ bar if RCMs are also included (Fig. 2). As better preservation/exposure could allow still smaller embedded craters to be uncovered, we interpret our fits as upper limits. The best fit to DTM1 (DTM2) ancient craters alone is 1.7 ± 0.3 bar (2.2 ± 0.3 bar), falling to 0.8 ± 0.1 bar (0.9 ± 0.1 mbar) if RCMs are included.

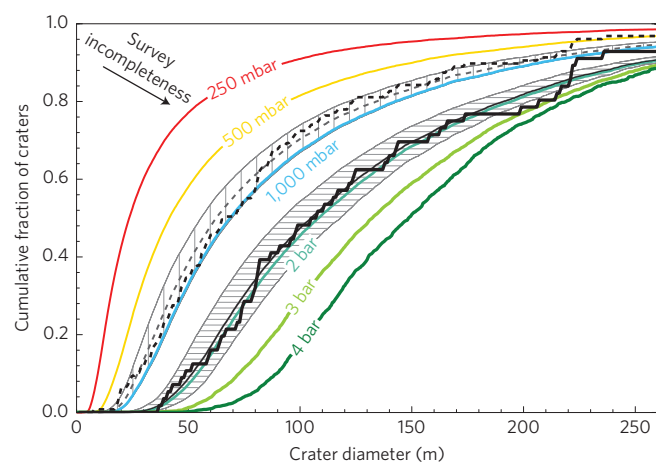


Figure 2 | Upper limits on the atmospheric pressure of early Mars.

Comparison of model crater size–frequency distributions to observations. Solid black line corresponds to definite embedded craters. Dashed black line additionally includes rimmed circular mesas. Stair-stepping in the data curves corresponds to individual craters. Coloured lines show model predictions for atmospheric filtering of small impactors at different pressures. Grey hatched regions correspond to 2σ statistical error envelopes around the best-fit palaeopressure to the data (best fits shown by thick grey lines). Survey incompleteness leads to overestimates of median crater size, so best fits are upper limits.

The results are sensitive to target strength, as expected²². Increasing the target rock-mass strength to a hard-rock-like 6.9 MPa (ref. 23) while holding all other parameters constant increases the combined (including RCMs) upper limit on P to ~ 2 bar (Supplementary Methods). Figure 2 assumes weak soil-like target strength appropriate for river alluvium in an aggrading sedimentary deposit: if sediment developed bedrock-like rock-mass strength by early diagenesis, the upper limit is greatly increased. Sensitivity tests show a relatively minor effect of fragmentation on the results (Supplementary Methods).

We do not consider crater shrinkage or expansion by crater degradation. Only shrinkage matters for the purpose of setting an upper bound on P : as the crater is abraded, the exposed radius must eventually vanish. We surmise that shrinkage is a small effect because impact craters are bowl-shaped (as opposed to cone-shaped) and because rims are frequently preserved.

Our technique rules out a thick, stable palaeoatmosphere and cannot exclude atmospheric collapse–re-inflation cycles on timescales much shorter than the sedimentary basin-filling time. General circulation models predict that atmospheric collapse to form CO_2 ice sheets and subsequent re-inflation might be triggered by obliquity change⁵. If sediment accumulated at $1\text{--}100\ \mu\text{m yr}^{-1}$ (ref. 20), our DTMs could integrate over $\sim 10^6\text{--}10^8$ years of sedimentation and contain many collapse–re-inflation cycles. Therefore one interpretation is that smaller ancient craters formed while the atmosphere was collapsed, whereas rivers formed during high-obliquity, thick-atmosphere intervals. However, published models indicate that collapse to form polar CO_2 ice sheets occurs only for pressures less than our upper limit⁵. If these models are correct, then our pressure constraint is a true upper bound on typical atmospheric pressure.

Downward revisions to the infrared opacity of CO_2 indicate that any amount of CO_2 is insufficient to warm early Mars \bar{T} to the freezing point⁵. Even if further work incorporating radiatively active clouds²⁴ moderates this conclusion, our result is an independent constraint on stable $\text{CO}_2/\text{H}_2\text{O}$ warm–wet solutions (Fig. 3). However, increased CO_2 below the warm–wet threshold primes the

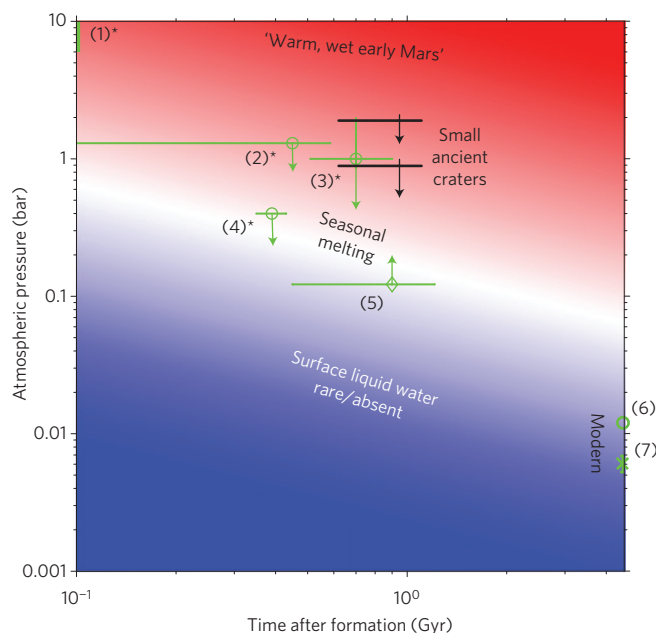


Figure 3 | Palaeopressure constraints on the great drying of Mars. Black symbols are the result from this work. Green symbols are other estimates (asterisks mark indirect estimates): (1*) cosmochemical estimate²⁶; (2*) prehnite stability; (3*) carbonate $\text{Mg}/\text{Ca}/\text{Fe}$ (ref. 27); (4*) $^{40}\text{Ar}/^{36}\text{Ar}$ (ref. 32); (5) bomb sag¹⁸; (6) modern atmosphere; (7) modern atmosphere + CO_2 ice. Approximate and model-dependent implications for sustained surface habitability are shown by background colours (blue, always below freezing; red, year-round melting; slope schematically shows the effect of the fainter young sun). ‘Warm, wet early Mars’ refers to the stable climate solution in ref. 2. Age estimates are detailed in Supplementary Methods.

climate of Mars for surface liquid water production by other relatively short-lived mechanisms, by adding to the greenhouse effect, pressure-broadening the absorption lines of other gases⁴, suppressing evaporitic cooling⁶ and increasing atmospheric heat capacity³.

If the small-crater limit is representative of early Mars P , it is difficult to envisage continuous stability of surface liquid water for the $10^4\text{--}10^5$ yr needed to allow water to cycle between deep aquifers and the surface. This is true even with optimistic CO_2 radiative forcing parameterizations. Transient warming by eruptions, impacts, or infrequent orbital conditions could unfreeze the surface and shallow subsurface, allowing runoff, but would not last long enough to unfreeze ground at ~ 1 km depth. Therefore, $\text{CO}_2/\text{H}_2\text{O}$ atmospheric models do not support \bar{T} above freezing on early Mars, which has implications for sedimentary rock formation and diagenesis, habitability and groundwater hydrology. (A recent study²⁵ shows that mean temperatures above the freezing point are marginally consistent with our result, but only if the early Mars atmosphere contained $\geq 20\%$ H_2).

Atmospheric loss must be part of the explanation for Mars’s great drying, if only because freshwater rivers cannot flow for hundreds of kilometres when simultaneously boiling and freezing. How high P was, and its decay over time, are not known. The 2014–2015 Mars Atmosphere and Volatile Evolution (MAVEN) mission will measure modern loss processes, which is complementary to our geologic palaeoproxy approach.

Mars would have formed with $\geq 6\text{--}10$ bars CO_2 equivalent of carbon assuming the same initial $[\text{C}]$ and $[\text{Kr}]$ as Earth. $^{40}\text{Ar}/^{36}\text{Ar}$ and $^{129}\text{Xe}/^{132}\text{Xe}$ suggest that 90–99% of the initial atmosphere was lost before ~ 4.1 Ga (ref. 26). Subsequent loss rates are less clear; Mars’s $\text{C}/^{84}\text{Kr}$ ratio suggests $P \sim 60$ mbar following the Late Heavy Bombardment²⁶.

Ref. 18 uses a volcanic bomb sag in Gusev crater to infer $P > 120$ mbar from the bomb sag's terminal velocity. This is consistent with our result. Our small-crater constraints on early Mars atmospheric pressure are also congruent with isotopic and mineralogic indicators, which generally require more assumptions than our method. For example, prehnite is observed on Mars and is unstable for CO_2 mixing ratios $> 2 \times 10^{-3}$. This implies $P \lesssim 1$ bar, but only if water at depth was in equilibrium with the atmosphere. The composition of a carbonate-rich outcrop at Gusev has been interpreted to require $P = 0.5\text{--}2$ bar assuming that the carbonates are a solid solution in thermodynamic equilibrium²⁷. Models¹⁶ of volcanic degassing, escape of CO_2 and impact delivery also hint that the atmospheric pressure on Mars at the time of the Late Heavy Bombardment was not greater than our estimate.

In the future, pyroclastic-blast runout length or even rainsplash²⁸ could be used to constrain P . Curiosity's field site in Gale crater contains syndepositional craters (Fig. 1a), so Curiosity could validate orbital identifications of embedded craters along its traverse. The 40-year-old prediction of a connection between drying and atmospheric decay could be tested by applying the small-crater technique to sedimentary deposits of different ages—ranging from Mawrth (the oldest known sedimentary sequence in the Solar System), through Meridiani, to the relatively young Valles Marineris silica deposits. This could yield a time series of constraints on the atmospheric pressure of early Mars, stratigraphically coordinated to the sedimentary record of Mars's great drying.

Methods

DTM generation. DTMs were constructed following ref. 11 from HiRISE images PSP_007474_1745/ESP_024497_1745 (DTM 1) and ESP_017548_1740/ESP_019104_1740 (DTM 2). Mars Orbiter Laser Altimeter Precision Experiment Data Records were used as ground control points. Optimal resolution depends on HiRISE image map scale (0.25–0.5 m), giving 1–2.5 m per post DTMs. Vertical precision is ~ 0.3 m, with 90% probability of precision $\lesssim 1$ m (Supplementary Methods).

Cratering model. We built a synthetic impactor population by drawing randomly from distributions of material types, densities and ablation coefficients, which are set based on terrestrial fireball network observations (the model assumes the same fractional distribution at Mars). Details of the cratering model, including the sources for parameter choices, are given in the Supplementary Methods. We advect these populations to the surface through atmospheres with scale height 10.7 km. The code does not track planet curvature; we do not allow impactors to skip back to space. The atmosphere drains kinetic energy from impactors through drag (per unit mass):

$$dv/dt = C_D \rho_a v^2 A$$

(we assume a drag coefficient $C_D = 1$ across the velocity (v) and size range of interest) and ablation:

$$dm/dt = (C_h \rho_a v^3 A)/2\zeta$$

where t is time, m is mass, ρ_a is local atmospheric density, A is cross-sectional area, C_h is the heat transfer coefficient and ζ is the heat of ablation. Particles braked to < 500 m s⁻¹ would not form hypervelocity craters and are removed from the simulation. We do not track secondary craters, because metre-sized endoatmospheric projectiles are likely to be braked to low speeds for the relatively thick atmospheres we are evaluating. In other words, if wet-era small craters are secondaries, then the atmosphere of early Mars was thin. Transient impact-induced increases in P would not affect our upper limit; transient local decreases in P could conceivably enhance secondary flux. It has been suggested that unrecognized secondary craters significantly contribute to all counts of $D < 1$ km martian craters²⁹ and, although it has been shown by others that unrecognized secondaries are not required to explain observed crater size–frequency distributions^{12,30,31}, the small-crater-rich size–frequency distribution of secondaries (if pervasively present with high relative frequency on all martian surfaces, as suggested by ref. 29) could mimic the effect of a lower atmospheric pressure. Crater sizes are calculated using π -group scaling²³, assuming a target strength of 65 kPa and a target density of 2,000 kg m⁻³ appropriate to cohesive desert alluvium¹⁴, with the Holsapple scaling parameters

$k_1 = 0.132$, $k_2 = 0.26$, $k_3 = 1.1$ and $\mu = 0.41$. We adopt the values in keith.aawashington.edu/craterdata/scaling/theory.pdf; note that the value $k_1 = 0.24$ given in Table 1 of ref. 23 is in error (K. A. Holsapple, personal communication).

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Author contributions

E.S.K. designed research, picked craters, carried out the data–model comparison and drafted the main text. J.-P.W. wrote the forward model of impactor–atmosphere interactions. A.L. built the digital terrain models and wrote the corresponding Supplementary text. O.A. supervised research. All authors contributed to the interpretation of the results and to the revisions.

Additional information

Supplementary information is available in the [online version of the paper](#). Reprints and permissions information is available online at www.nature.com/reprints. Correspondence and requests for materials should be addressed to E.S.K.

Competing financial interests

The authors declare no competing financial interests.